Symmetry saga on LHC

Vladimir Shevchenko

National Research Centre «Kurchatov Institute»

Jam Jamor 2008

References and acknowledgements

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http://cern.ch

and sites of collaborations. The speaker is thankful to A.Golutvin, T.Nakada, A.Vorobyev, A.Bondar, V.Egorychev, V.Belyaev, D.Kharzeev, M.Polikarpov, V.Zakharov and many other colleagues for numerous discussions of LHC physics. Slides from conference talks and public lectures of A.Golutvin, R.Forty, D.Kharzeev, B.Schenke, S.Myers are borrowed.

PLAN:

1. Theory

- Scales, matter and dynamics where we are
- Main ideas of the Standard Model
- Symmetries and their violation

2. Experiment

- ABC of LHC
- Four LHC experiments from symmetry point of view
- Recent results and prospectives



«Cube of theories» (M.Bronshtein, A. Zelmanov)



Discovery and formulation of the Standard Model is perhaps the most outstanding result of human intellect for all history of mankind



event of the same decade.»

R.Feynman

What are the key pillars of the Standard Model?

The Standard Model Zoo

Gauge group SU(3) ◊ SU(2)_L ◊ U(1)_Y [g; W, Z; ^y_o]

is broken down to $SU(3) \otimes U(1)_Q$

Quarks and leptons form *generations*. Masses come out of *interactions*



Mass hierarchy (from hep-ph/0603118)

Elementary Particles



Matter: where from?

Elementary











Dynamics: where from? From symmetry and its breaking

Main ideas of the Standard Model

1. Gauge invariance

There is an internal symmetry in the theory with respect to some special field transformations. Different configurations of dynamical variables describe one and the same set of physical observables.

Theory is local (and hence elegant and economical) in terms of non-observable potentials, while non-local (and hence cumbersome) in terms of observable field strengths)

 $F_{\mu\nu}(x)$ non-local interaction $A_{\mu}(x)$ local interaction with currents with currents

2. Spontaneous symmetry breaking

Symmetry of equation is not always a symmetry of its solution. If solution is to realize extremum of some quantity, it can be reached on less symmetric configuration.





Y.Nambu

Elementary example: opaque square



 $2\sqrt{2} \approx 2.82$ $1 + \sqrt{3} \approx 2.73$ $\sqrt{2} + \sqrt{\frac{3}{2}} \approx 2.64$

3. Masses come from broken symmetry

Ginzburg-Landau-Higgs mechanism

V.L.Ginzburg, L.D.Landau, 1950 P.W.Higgs, 1964; F.Englert, R.Brout, 1964 G.S.Guralnik, C.R.Hagen, T.W.Kibble, 1964

as if gauge symmetry is "broken" $SU(2)_{L} \otimes U(1)_{Y} \rightarrow U(1)_{Q}$

$$L = |(i\partial + A)\Phi|^2 + V(\Phi)$$

- Electroweak sector of SM
- Superconductivity
- Meissner scenario for confinement

essentially classical

Dimensional transmutation

G. 't Hooft, 1972 D.Gross, F.Wilczek, 1973 H.Politzer, 1973 scale invariance is broken by anomaly SU(3)

$$\Lambda = \mu \exp\left(-\frac{2\pi}{b\alpha_s(\mu)}\right)$$

$$T^{\mu}_{\mu} = \frac{\beta(g)}{2g^3} \operatorname{Tr} F_{\mu\nu} F^{\mu\nu}$$

principally quantum

Standard Model has many built in small parameters. Each one gives a chance to construct perturbation theory. Some are related to symmetry breaking.

- 1. Dimensionless interaction constants
- 2. Rank of color gauge group SU(3)
- 3. Quark masses in strong interaction scale units
- 4. Strong to weak scales ratio
- 5. Yukawa constants
- 6. Quark mixing parameters

 $1/9 = 1/3^2$ $\frac{m_{u,d}}{\Lambda} = (0.5 \div 2)\%$ $G_F m_p^2 = 1 \cdot 10^{-5}$ $\frac{m_e}{m_t} = 3 \cdot 10^{-6}$

 $\alpha_{em}^{-1} = 137$ $\alpha_s^{-1}(M_Z) = 8.5$

$$\lambda = |V_{us}| = 0.22$$

7. ...

Many important small parameters are beyond the Standard Model:

$$m_{\nu}$$
 $G_N m_p^2 = 6 \cdot 10^{-39}$



Basic classification of symmetries goes along these lines: type of transformation (e.g. discrete or continuous) and accuracy of comparison (built in or observer-dependent)

Maximal symmetry is not very interesting..

true harmony = beauty **slightly** broken symmetry





But not too much broken...



Besides continuous symmetries (local and global) very important role is played by discrete ones. Among them the most important are:

- $P parity transformation X \rightarrow -X$ det P = -1!
- T time reversal $t \rightarrow -t$
- **C** charge conjugation $e \rightarrow -e$

In classical physics

In quantum field theory the status of transformations C, P, T or their pairs CP, CT, TP is different from **CPT**.

Namely, for any local Lorentz-invariant quantum field theory one has

CPT theorem

J.Schwinger, 1951 G.Lüders, W.Pauli, 1954

Physically, **CPT** theorem means that antiparticles and their interactions are indistinguishable from particles, going along the same world-lines in 3+1 space-time but in opposite directions.

In particular, mass of any particle regardless of its dynamics exactly equals to the mass of its antiparticle (experimentally checked for K-mesons at the level 1 to 10^{18}).

C, P, T have no such kinematical status

P-parity - broken is macroscopic world





Our world

DNA

Wonderland

P-parity for elementary particles



T.D.Lee, C.N.Yang, 1956



C.S.Wu, 1957



In Standard Model P-parity violation is built in at the level of Lagrangian

Thus V-A effective Lagrangian looks CP-invariant











L.D.Landau, 1957: CP-parity conservation hypothesis J.Cronin, V.Fitch, 1964: Discovery of CP-nonconservation in neutral kaons decays

CPLEAR 1999





M. Kobayashi, T.Maskawa, 1974: theoretical mechanism of

CP-violation

based on ideas of *N.Cabibbo, 1963*



Idea: flavor eigenstates are superposition of mass eigenstates (and vice versa)

In other way: it is impossible to diagonalize interaction term for charged currents and free kinetic term

$$L_{\text{int}} = \frac{g_2}{\sqrt{2}} \langle \overline{q}_L, \overline{c}_L, \overline{t}_L \rangle \gamma^{\mu} \hat{V}_{CKM} \begin{pmatrix} d_L \\ s_L \\ b_L \end{pmatrix} W^+_{\mu} + h.c.$$

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

It can be shown that arbitrary complex unitary N N matrix can be parameterized by N(N-1)/2 Euler angles and (N-1)(N-2)/2 complex phases.

For N<3 any matrix can be unitary rotated to equivalent real.

On the other hand, there are such 3 3 matrices that no choice of quark fields phases can make them real.

$$\eta = \arg\left(\frac{V_{ij}V_{ik}^*}{V_{lj}V_{lk}^*}\right) \qquad J = \left|\operatorname{Im} V_{i\alpha}V_{j\beta}V_{i\beta}^*V_{j\alpha}^*\right| \sim 3 \times 10^{-5}$$

Remark to remember:

is'n it strange that strong interactions do not violate **CP** at Lagrangian level?

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - \frac{n_f g^2 \theta}{32\pi^2}F_{\mu\nu}\tilde{F}^{\mu\nu} + \bar{\psi}(i\gamma^{\mu}D_{\mu} - me^{i\theta'\gamma_5})\psi$$

Current experimental constraint on θ is $|\theta| \le 10^{-9}$

According to **CPT** theorem **CP**-violation means **T**-violation. The problem of time arrow (*A.Eddington, 1927*)



Important general conclusion



SYMMETRY



can be broken at the level of «laws» DYNAMICS AS SYMMETRY BREAKING

or can be broken spontaneously, at the level of «states»



Particle accelerators used in industry may be thought of as sophisticated hammers. But accelerators used in basic science (and Large Hadron Collider at CERN in particular) are better to think of as microscopes, not hammers.



We need the highest possible resolution to test the fundamental symmetries of our Universe

Jet d'Eau 140 m

Mont Blanc, 4808 m



CERN (Conseil Européen pour la Recherche Nucléaire) founded 29 September 1954

CMS



ABC of LHC

- Tonnel length 27 kilometers
- Tonnel depth between 50 and 175 meters
- Colliding p-p beams, 2808 bunches, 1.15 10¹¹ p/bunch
- v = 0.9999998 c
- Energy = 7 TeV + 7 TeV, luminosity ~ 10^{34} cm⁻² sec⁻¹
- About 600 millions collisions per second

Proton energy = $7 \text{ TeV} = 10^{-6} \text{ joul.}$

This is about kinetic energy of a fly:



Why not to use flies in particle physics? Besause Avogadro number = 6.022×10^{23} (mol)⁻¹ Fly's energy is distributed among ~ 10^{22} nucleons. On the other hand, total beam energy is given by 2808 bunches × 10^{11} protons/bunch × 7 TeV/proton ~ 400 MJ This is about 100 kg TNT-equivalent or kinetic energy of «Admiral Kuznetsov» cruising at 8 knots.



Modern HEP experiment has a long way to go ...

The life of an experiment

- 1984 Workshop in Lausanne on installing a Large Hadron Collider (LHC) in the LEP tunnel
- 1987 CERN's long-range planning committee chaired by Carlo Rubbia recommends LHC as the right choice for lab's future
- **1989** ECFA Study Week on instrumentation technology for a high-luminosity hadron collider; Barcelona; LEP collider starts operation
- 1990 ECFA LHC workshop, Aachen
- 1992 General meeting on LHC physics and detectors, Evian-les-Bains
- 1993 Letters of intent for LHC detectors submitted
- 1994 Technical proposals for ATLAS and CMS approved
- 1998 Construction begins
- 2000 CMS assembly begins above ground; LEP collider closes
- 2003 ATLAS underground cavern completed and assembly started
- 2004 CMS cavern completed
- 2007 Experiments ready for beam
- 2007 First proton-proton collisions
- 2008 First results
- 2010 Reach design luminosity
- >2014 Upgrade LHC luminosity by factor of 10





ECFA 84/85 CERN 84-10 5 September 198











Two beam pipes with common magnetic system

Magnets

Dipole superconducting **Ti-Nb** magnets, magnetic field 8.3 T, current 11700 A, temperature 1.9 K (700000 litres of liquid **He**)





1232 items, 0.5M CHF each

r [m] B [T] = 3.33 p [GeV]27 km, 2/3 is covered by dipole magnets field $r \sim 2800$ m B = 8.3 T p = 7 TeV

Energy stored in magnets ~ 10.4 GJ = kinetic energy of Boeing at 700 km/h



Testing symmetries in LHC experiments



General purpose - everything with high enough p_T

Electroweak gauge symmetry breaking pattern: Higgs boson or New Physics?

Space-time symmetries: extra dimensions, black holes, KK-states?

Supersymmetry: particles – superpartners? Dark matter?





Enigma of flavor

CP-symmetry violation: new sources? Baryon asymmetry of the Universe. Indirect search of superpartners.



New state of matter - new symmetries?

Chiral symmetry of strong interactions: pattern of restoration? Deconfinement. P-parity violation as interplay of strong and electromagnetic interactions?

ATLAS – A Toroidal Lhc ApparatuS The largest detector in history of particle physics – 45 meters and 7000 tons



CMS – Compact Muon Spectrometer

21 meters и 12500 tons



















Brief History of the Standard Model



Hunting for Higgs...

ynn





 $(\phi)^{\dagger}D^{*}\phi - U(\phi) - \frac{1}{4}F_{\mu\nu}F^{\mu\nu}F^{\mu\nu}$ 5t+ + B (++)2





		ATLAS Exotics Searches* - 95% CL Lower Limits (Status: Dec. 2011)		
			гтт	
	Large ED (ADD) : monojet	L=1.0 fb ⁻¹ (2011) [ATLAS-CONF-2011-096] 3.2 TeV $M_D(\delta=2)$		
	Large ED (ADD) : diphoton	L=2.1 fb ⁻¹ (2011) [Preliminary] 3.0 TeV M_S (GRW cut-off) ATLA	AS	
	$UED: \gamma\gamma + E_{T,miss}$	L=1.1 fb ⁻¹ (2011) [arXiv:1111.4116] 1.23 TeV Compact. scale 1/R (SPS8) Prelimit	nary	
ra dimensions	RS with $k/M_{Pl} = 0.1$: $\gamma\gamma$, ee, $\mu\mu$ combined, $m_{\gamma\gamma, \parallel}$	L=1.1-2.1 fb ⁻¹ (2011) [Preliminary, arXiv:1108.1582] 1.95 TeV Graviton mass	1	
	RS with $k/M_{Pl} = 0.1$: ZZ resonance, m_{IIII}	$L=1.0 \text{ fb}^{-1} (2011) \text{ [ATLAS-CONF-2011-144]} 575 \text{ GeV} \text{ Graviton mass} \qquad \int L dt = (0.03 - 2.1)$	fb ⁻ '	
	RS with g_{aqgKK}/g_{s} = -0.20 : $H_{T} + E_{T,miss}$	L=1.0 fb ⁻¹ (2011) [ATLAS-CONF-2011-123] 840 GeV KK gluon mass	TeV	
	Quantum black hole (QBH) : m_{dijet} , $F(\chi)$	L=36 pb ⁻¹ (2010) [arXiv:1103.3864] 3.67 TeV M_D (δ =6)		
Ext	QBH : High-mass $\sigma_{t + X}$	L=33 pb ⁻¹ (2010) [ATLAS-CONF-2011-070] 2.35 TeV M _D		
	ADD BH ($M_{TH}/M_{D}=3$) : multijet, Σp_{T} , N_{jets}	L=35 pb ⁻¹ (2010) [ATLAS-CONF-2011-068] 1.37 TeV $M_{\rm D}~(\delta=6)$		
	ADD BH ($M_{TH}/M_{D}=3$) : SS dimuon, $N_{ch. part.}$	L=1.3 fb ⁻¹ (2011) [arXiv:1111.0080] 1.25 TeV M _D (δ=6)		
	ADD BH ($M_{TH}/M_{D}=3$) : leptons + jets, Σp_{T}	L=1.0 fb ⁻¹ (2011) [ATLAS-CONF-2011-147] 1.5 TeV M_D (δ =6)		
~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~	qqqq contact interaction : $F_{\chi}(m_{\text{dijet}})$	L=36 pb ⁻¹ (2010) [arXiv:1103.3864 (Bayesian limit)] 6.7 TeV		
0	qqll contact interaction : ee, $\mu\mu$ combined, $m_{\mu}$	L=1.1-1.2 fb ⁻¹ (2011) [Preliminary] 10.2 TeV Λ (constructive int.)		
2	SSM : m _{ee/µµ}	L=1.1-1.2 fb ⁻¹ (2011) [arXiv:1108.1582] 1.83 TeV Z' mass		
	SSM : $m_{T,e/\mu}$	L=1.0 fb ⁻¹ (2011) [arXiv:1108.1316] 2.15 TeV W' mass		
3	Scalar LQ pairs ( $\beta$ =1) : kin. vars. in eejj, evjj	L=1.0 fb ⁻¹ (2011) [Preliminary] 660 GeV 1 st gen. LQ mass		
L(	Scalar LQ pairs ( $\beta$ =1) : kin. vars. in µµjj, µvjj	L=35 pb ⁻¹ (2010) [arXiv:1104.4481] 422 GeV 2 nd gen. LQ mass		
- Ue	$4^{th}$ generation : coll. mass in Q $\overline{Q}_{4} \rightarrow WqWq$	L=37 pb ⁻¹ (2010) [CONF-2011-022] 270 GeV Q ₄ mass		
h g	$4^{\text{th}}$ generation : $d_{a} \overline{d}_{a} \rightarrow Wt \tilde{W}t$ (2-lep SS)	L=34 pb ⁻¹ (2010) [1108.0366] 290 GeV d ₄ mass		
4-1	$TT_{exp, 4th gen} \rightarrow t\bar{t} + A_0 A_0^2$ : 1-lep + jets + $E_{T, miss}$	L=1.0 fb ⁻¹ (2011) [arXiv:1109.4725] 420 GeV T mass ( $m(A_{o}) < 140$ GeV)		
	Techni-hadrons : dilepton, mee/uu	L=1.1-1.2 fb ⁻¹ (2011) [CONF-2011-125] 470 GeV $\rho_{-}/\omega_{T}$ mass $(m(\rho_{-}/\omega_{T}) - m(\pi_{T}) = 100 \text{ GeV})$		
	Major. neutr. (LRSM, no mixing) : 2-lep + jets	L=34 pb ⁻¹ (2010) [ATLAS-CONF-2011-115] 780 GeV N mass $(m(W_p) = 1 \text{ TeV})$		
	Major. neutr. (LRSM, no mixing) : 2-lep + jets	L=34 pb ⁻¹ (2010) [ATLAS-CONF-2011-115] 1.350 TeV W _B mass (230 < m(N) < 700 GeV)		
	$H_{L}^{\pm\pm}$ (DY prod., BR( $H^{\pm\pm} \rightarrow \mu\mu$ )=1) : $m_{\mu\nu}$ (like size)	<i>L</i> =1.6 fb ⁻¹ (2011) [CONF-2011-127] 375 GeV $H_{1}^{\pm\pm}$ mass		
er	Excited quarks : $\gamma$ -jet resonance, $m_{\text{vist}}^{\mu\mu}$	L=2.1 fb ⁻¹ (2011) [Preliminary] 2.46 TeV q [*] mass		
Oth	Excited quarks : dijet resonance, $m_{\text{dijet}}^{\eta \text{or}}$	L=1.0 fb ⁻¹ (2011) [arXiv:1108.6311] 2.99 TeV q [*] mass		
	Axigluons : m _{dijet}	L=1.0 fb ⁻¹ (2011) [arXiv:1108.6311] 3.32 TeV Axigluon mass		
	Color octet scalar : $m_{\text{dijet}}$	L=1.0 fb ⁻¹ (2011) [arXiv:1108.6311] 1.92 TeV Scalar resonance mass		
	Vector-like quark : CC, m	L=1.0 fb ⁻¹ (2011) [Preliminary] 900 GeV Q mass (coupling $\kappa_{oO} = v/m_O$ )		
	Vector-like quark : NC, $m_{lig}$ $l_{=1.0 \text{ fb}^{-1}}$ (2011) [Preliminary] 760 GeV Q mass (coupling $\kappa_{roc} = v/m_o$ )			
		10 ⁻¹ 1 10	10 ²	
		Maga apple	$T_{-} $	

*Only a selection of the available results leading to mass limits shown

Mass scale [TeV]

ATLAS SUSY Searches* - 95% CL Lower Limits (Status: Dec. 2011)
MSUGRA/CMSSM : 0-lep + j's + $E_{T,miss}$ L=1.0 fb ⁻¹ (2011) [arXiv:1109.6572] 950 GeV $\tilde{q} = \tilde{g}$ mass ATLAS
$MSUGRA/CMSSM : 1-lep + j's + E_{T,miss} = L_{\pm 1.0 \text{ fb}^{-1}(2011) \text{ [arXiv:1109.6606]}} 820 \text{ GeV}  \tilde{q} = \tilde{g} \text{ mass} \qquad Preliminary$
$MSUGRA/CMSSM : multijets + E_{T,miss} \xrightarrow{L=1.3 \text{ fb}^{-1}(2011) [arXiv:1110.2299]} 680 \text{ GeV}  \widetilde{g} \text{ mass } (\text{for } m(\widetilde{q}) = 2m(\widetilde{g}))$
Simpl. mod. : 0-lep + j's + $E_{T,miss}$ L=1.0 fb ⁻¹ (2011) [arXiv:1109.6572] 1.075 TeV $\tilde{q} = \tilde{g}$ mass (light $\tilde{\chi}_1^0$ ) $\int Ldt = (0.03 - 2.0)$ fb ⁻¹
Simpl. mod. : 0-lep + j's + $E_{T,miss}$ L=1.0 fb ⁻¹ (2011) [arXiv:1109.6572] 875 GeV $\tilde{q}$ mass $(m(\tilde{g}) < 2$ TeV, light $\tilde{\chi}_1^0$ ) Is = 7 TeV
Simpl. mod. : 0-lep + j's + $E_{T,miss}$ L=1.0 fb ⁻¹ (2011) [arXiv:1109.6572] 700 GeV $\tilde{g}$ mass $(m(\tilde{q}) < 2 \text{ TeV}, \text{ light } \tilde{\chi}_1^0)$
Simpl. mod. : 0-lep + j's + $E_{T,miss}$ L=1.0 fb ⁻¹ (2011) [ATLAS-CONF-2011-155] 700 GeV $\tilde{q}$ mass $(m(\tilde{g}) < 2 \text{ TeV}, m(\tilde{\chi}_1^0) < 200 \text{ GeV})$
Simpl. mod. : 0-lep + j's + $E_{T,miss}$ L=1.0 fb ⁻¹ (2011) [ATLAS-CONF-2011-155] 650 GeV $\tilde{g}$ mass $(m(\tilde{q}) < 2 \text{ TeV}, m(\tilde{\chi}_1^0) < 200 \text{ GeV})$
Simpl. mod. $(\widetilde{g} \rightarrow q \overline{q} \widetilde{\chi}^{\pm})$ : 1-lep + j's + $E_{T,miss}$ L=1.0 fb ⁻¹ (2011) [arXiv:1109.6606] 600 GeV $\widetilde{g}$ mass $(m(\widetilde{\chi}^0_1) < 200 \text{ GeV}, \Delta m(\widetilde{\chi}^{\pm}, \widetilde{\chi}^0) / \Delta m(\widetilde{g}, \widetilde{\chi}^0) > 1/2)$
Simpl. mod. : 0-lep + b-jets + j's + $E_{T,miss}$ L=0.83 (b ⁻¹ (2011) [ATLAS-CONF-2011-098] 720 GeV $\tilde{g}$ mass ( $m(\tilde{b})$ < 600 GeV, light $\tilde{\chi}_{1}^{0}$ )
Simpl. mod. $(\tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0)$ : 1-lep + b-jets + j's + $E_{T,miss}$ L=1.0 fb ⁻¹ (2011) [ATLAS-CONF-2011-130] 540 GeV $\tilde{g}$ mass $(m(\tilde{\chi}_1^0) < 80 \text{ GeV})$
Simpl. mod. $(\tilde{b}_1 \rightarrow b \tilde{\chi}_1^0)$ : 2 b-jets + $E_{T,miss}$ L=2.05 fb ⁻¹ (2011) [Preliminary] 390 GeV $\tilde{b}$ mass $(m(\tilde{\chi}_1^0) < 60 \text{ GeV})$
Simpl. mod. $(\tilde{\chi}_{1}^{\pm}\tilde{\chi}_{2}^{0} \rightarrow 3   \tilde{\chi}_{1}^{0})$ : 2-lep SS + $E_{T,\text{miss}}$ $L=1.0 \text{ fb}^{-1}$ (2011) [arXiv:1110.6189] $\tilde{\chi}_{1}^{\pm}$ mass (light $\tilde{\chi}_{1}^{0}, m(\tilde{l}) = \frac{1}{2}(m(\tilde{\chi}_{1}^{\pm}) + m(\tilde{\chi}_{2}^{0})))$
GMSB : 2-lep OS _{SF} + $E_{T,miss}$ L=1.0 fb ⁻¹ (2011) [ATLAS-CONF-2011-156] 810 GeV $\tilde{g}$ mass (corresp. to $\Lambda < 35$ TeV, tan $\beta < 35$ )
GGM + Simpl. model : $\gamma\gamma + E_{T,miss}$ L=1.1 fb ⁻¹ (2011) [arXiv:1111.4116] 805 GeV $\tilde{g}$ mass (m(bino) > 50 GeV)
GMSB : stable $\tilde{\tau}$ L=37 pb ⁻¹ (2010) [1106.4495] $\tilde{\tau}$ mass
AMSB : long-lived $\tilde{\chi}_{1}^{\pm}$ L=1.0 fb ⁻¹ (2011) [Prel] $\tilde{\chi}_{1}^{\pm}$ mass (0.5 < $\tau(\tilde{\chi}_{1}^{\pm})$ < 2 ns )
Stable massive particles : R-hadrons L=34 pb ⁻¹ (2010) [arXiv:1103.1984] 562 GeV g mass
Stable massive particles : R-hadrons L=34 pb ⁻¹ (2010) [arXiv:1103.1984] 294 GeV b mass
Stable massive particles : R-hadrons L=34 pb ⁻¹ (2010) [arXiv:1103.1984] 309 GeV t mass
Hypercolour scalar gluons : 4 jets, $m_{ij} \approx m_{kl}$ $L=34 \text{ pb}^{-1} (2010) [arXiv:1110.2693]$ sgluon mass (excl: $m_{sq} < 100 \text{ GeV}, m_{sq} \approx 140 \pm 3 \text{ GeV}$ )
RPV : high-mass eµ $L_{=1.1 \text{ fb}^{-1} (2011) [arXiv:1109.3089]}$ 1.32 TeV $\tilde{V}_{\tau}$ mass ( $\lambda_{311}^2 = 0.10, \lambda_{312}^2 = 0.05$ )
Bilinear RPV : 1-lep + j's + $E_{T,miss}$ $L=1.0 \text{ fb}^{-1}(2011) [arXiv:1109.6606]$ 760 GeV $\tilde{q} = \tilde{g} \text{ mass} (c\tau_{LSP} < 15 \text{ mm})$
10 ⁻¹ 1 10

*Only a selection of the available results leading to mass limits shown

Mass scale [TeV]

# LHCb Event Display







#### Bound states of **b** and anti-**b** quarks





# The primary goal of ALICE experiment is to study new state of matter, sometimes referred to as quark-gluon plasma It is not plasma







Coordinate space asymmetry is transformed into momentum space asymmetry

$$E\frac{d^{3}N}{d^{3}p} = \frac{1}{2\pi} \frac{d^{2}N}{p_{t}dp_{t}dy} \left(1 + \sum_{n=1}^{\infty} 2v_{n} \cos[n(\phi - \Psi_{R})]\right)$$

Single event – no special symmetry

 $1+2v_1\cos(\phi-\Psi_1)+2v_2\cos(2(\phi-\Psi_2))+2v_3\cos(3(\phi-\Psi_3))+$  $+2v_4\cos(4(\phi-\Psi_4))+2v_5\cos(5(\phi-\Psi_5))+...$ 



The matter produced at LHC still behaves as very low viscosity fluid

(from PRL, 105 (2010) 252302)

# Hydrodynamic description: $T^{\mu\nu} = (\epsilon + P)u^{\mu}u^{\nu} - Pg^{\mu\nu} + \pi^{\mu\nu}$ $\pi^{\mu\nu} = \pi^{\mu\nu}_{(1)} = \eta \left( \nabla^{\mu} u^{\nu} + \nabla^{\nu} u^{\mu} - \frac{2}{3} \Delta^{\mu\nu} \nabla_{\alpha} u^{\alpha} \right)$ Equation of motion: $\partial_{\mu} T^{\mu\nu} = 0$ Equation of state: $\epsilon = \epsilon(P)$

Emergent conformal symmetry for effective theory:

$$T^{\mu}_{\mu} = 0$$





Elliptic flow does not change much from RHIC to LHC

(from PRL, 105 (2010) 252302)

#### P-violation in heavy ions collisions?



# Comparison of magnetic fields



The Earths magnetic field	0.6 Gauss		
A common, hand-held magnet	100 Gauss		
The strongest steady magnetic fields achieved so far in the laboratory	4.5 x 10 ⁵ Gauss		
The strongest man-made fields ever achieved, if only briefly	10 ⁷ Gauss		
Typical surface, polar magnetic fields of radio pulsars	10 ¹³ Gauss		
Surface field of Magnetars	10 ¹⁵ Gauss		
http://solomon.as.utexas.edu/~duncan/magnetar.html			



Heavy ion collisions: the strongest magnetic field ever achieved in the laboratory Off central Gold-Gold Collisions at 100 GeV per nucleon  $e B(\tau=0.2 \text{ fm}) = 10^3 \sim 10^4 \text{ MeV}^2 \sim 10^{17} \text{ Gauss}$ 

(slide from D.Kharzeev)

A seminal suggestion: Kharzeev, Pisarski, Tytgat, '98; Halperin, Zhitnitsky, '98; Kharzeev, '04; Kharzeev, McLerran, Warringa '07; Kharzeev, Fukushima, Warringa '08

Possible experimental manifestations of

#### chiral magnetic effect

in heavy ion collisions ?



 $\mathbf{j} = \frac{e^2}{2\pi^2} \mu_5 \mathbf{B}$ 

Many complementary ways to derive (Chern-Simons, linear response, triangle loop etc). At effective Lagrangian level

$$\mu_5 \sim \dot{\theta}$$

Robust theoretical effect



P-parity forbidden correlators may be nonzero in external fields. Manifestation of chiral magnetic effect?

$$i \int dx \ e^{iq(x-y)} \left\langle \Omega | j_{\mu}(x) \cdot \partial j^{5,a}(y) | \Omega \right\rangle = \operatorname{Tr} \left[ Q^2 t^a \right] \cdot \left( -\frac{N_c}{4\pi^2} \right) \cdot q_{\nu} \tilde{F}_{\mu\nu}$$

Subtle interplay of strong and electromagnetic anomalies

# Conclusion

1). By the end of 2012 Higgs boson of the Standard Model will be either found (with the mass about 128 GeV) or ruled out in ATLAS and CMS experiments

2). In the latter case new effects in interactions of W and Z bosons are to be seen

3). By the end of 2012 a few minimal New Physics scenarios will be closed (or their effects seen)

4). Only SM Higgs in ATLAS and CMS and nothing in LHCb – Nightmare Scenario

THERE ARE MORE THINGS IN HEAVEN AND EARTH, HORATIO, THAN ARE DREAMT OF IN YOUR PHILOSOPHY.

W.SHAKESPEARE, HAMLET ACT 1, SCENE 5 Все, что на волю Высшую согласно, Своею волей чуждую творит, И под личиной вещества бесстрастной Везде огонь Божественный горит.

В.С.Соловьев